

# Appendix I

## Fixed- $(m, M)$ occupation numbers

Our purpose here is to derive a simple expression for the occupation probabilities  $\langle n_{m_{z_i}} \rangle^{mM}$  for  $m$  fermions in  $N$  sp states labeled by  $J_z$  quantum number  $m_{z_i}$ . Here,  $M$  are the eigenvalues of the  $J_z$  operator. As  $\langle n_{m_{z_i}} \rangle^{mM}$  is an expectation value, we can write a polynomial expansion in terms of the  $J_z$  operator [Dr-77],

$$\langle n_{m_{z_i}} \rangle^{mM} = \sum_{\mu} \langle n_{m_{z_i}} P_{\mu}(\hat{J}_z) \rangle^m P_{\mu}(\widehat{M}), \quad (I1)$$

where  $\widehat{M} = M/\sigma_{J_z}(m)$ ,  $\hat{J} = J_z/\sigma_{J_z}(m)$  and  $P_{\mu}(M)$  are orthogonal polynomials defined by the density  $\rho_{J_z}(M)$  which is close to a Gaussian. Retaining terms up to order 2, the expansion is,

$$\begin{aligned} \langle n_{m_{z_i}} \rangle^{mM} &= \langle n_{m_{z_i}} \rangle^m + \langle n_{m_{z_i}} \hat{J}_z \rangle^m \widehat{M} + \frac{(\langle n_{m_{z_i}} \hat{J}_z^2 \rangle^m - \langle n_{m_{z_i}} \rangle^m)(\widehat{M}^2 - 1)}{\langle J_z^4 \rangle^m - 1} \\ &= \langle n_{m_{z_i}} \rangle^m + \frac{\langle n_{m_{z_i}} J_z \rangle^m M}{\sigma_{J_z}^2(m)} \\ &+ \frac{1}{2} \left( \frac{\langle n_{m_{z_i}} J_z^2 \rangle^m}{\sigma_{J_z}^2(m)} - \langle n_{m_{z_i}} \rangle^m \right) \left( \frac{M^2}{\sigma_{J_z}^2(m)} - 1 \right). \end{aligned} \quad (I2)$$

In the above expression we used  $\langle \hat{J}_z^4 \rangle^m = 3$ , the value for a Gaussian  $\rho_{J_z}(M)$ . Now the formulas for the traces in Eq. (I2) are as follows. Firstly,

$$\langle n_{m_{z_i}} \rangle^m = \frac{m}{N} \langle \langle n_{m_{z_i}} \rangle \rangle^1 = \frac{m}{N}. \quad (\text{I3})$$

This implies  $n_{m_{z_i}}^{v=0} = \hat{n}/N$ . Also,

$$\langle n_{m_{z_i}} J_z \rangle^m = \frac{m(N-m)}{N(N-1)} \langle \langle n_{m_{z_i}} J_z \rangle \rangle^1, \quad (\text{I4})$$

with  $\langle \langle n_{m_{z_i}} J_z \rangle \rangle^1 = m_{z_i}$ . Unitary decomposition of the number operator gives,

$$\langle n_{m_{z_i}} J_z^2 \rangle^m = \langle n_{m_{z_i}} \rangle^m \langle J_z^2 \rangle^m + \langle n_{m_{z_i}}^{v=1} J_z^2 \rangle^m, \quad (\text{I5})$$

$$\langle n_{m_{z_i}}^{v=1} J_z^2 \rangle^m = \frac{m(N-m)(N-2m)}{N(N-1)(N-2)} \langle \langle n_{m_{z_i}}^{v=1} J_z^2 \rangle \rangle^1. \quad (\text{I6})$$

Now,

$$\langle \langle n_{m_{z_i}}^{v=1} J_z^2 \rangle \rangle^1 = \langle \langle (n_{m_{z_i}} - n_{m_{z_i}}^{v=0}) J_z^2 \rangle \rangle^1 = m_{z_i}^2 - \frac{1}{N} \langle \langle J_z^2 \rangle \rangle^1, \quad (\text{I7})$$

where we have used the result that  $n_{m_{z_i}}^{v=0} = \hat{n}/N$  deduced from Eq. (I3). Thus,

$$\langle n_{m_{z_i}} J_z^2 \rangle^m / \sigma_{J_z}^2(m) = \langle n_{m_{z_i}} \rangle^m + \frac{(N-2m)(m_{z_i}^2 - \langle J_z^2 \rangle^1)}{(N-2)N \langle J_z^2 \rangle^1}. \quad (\text{I8})$$

Substituting above traces in Eq. (I2) we have,

$$\langle n_{m_{z_i}} \rangle^{m,M} = \bar{m} + \frac{m_{z_i} M}{N \langle J_z^2 \rangle^1} - \frac{(\bar{m} - 1/2)(m_{z_i}^2 - \langle J_z^2 \rangle^1)(M^2 - \sigma_{J_z}^2(m))}{N^2 (\langle J_z^2 \rangle^1)^2 \bar{m}(1 - \bar{m})}, \quad (\text{I9})$$

where  $\bar{m} = m/N$ . The expression for the occupation number  $\langle n_{m_{z_i}} \rangle^{m,M}$  is close to that obtained in [Mu-00, Ze-04] where statistical mechanics approach has been employed. Thus, we have successfully reproduced the previously obtained results using moment method formalism.